\& Rollett is really significant, that a considerable strain in a bond angle may be tolerated if a more satisfactory packing or hydrogen bond pattern may be so achieved. The individual molecules in the crystal are held together by a three-dimensional network of hydrogen bonds in a manner which might well serve as a model to demonstrate the importance and effectiveness of such bond formation in determining the molecular packing within a crystal. The thoroughness and simplicity of the scheme may be seen from Fig. 3.

The two hydrogen atoms of the amino group both take part in bond formation, in the one case to form an intra-molecular linkage $\mathrm{N}(1)-\mathrm{H} \cdots \mathrm{O}\left(\mathrm{l}^{\prime}\right)$, in the other the bond $\mathrm{N}(1)-\mathrm{H} \cdots \mathrm{O}\left(\mathrm{H}_{2} \mathrm{O}\right)$. In turn, the two hydrogens of the water molecule participate in the bonds $\mathrm{O}\left(\mathrm{H}_{2} \mathrm{O}\right)-\mathrm{H} \cdots \mathrm{O}(2), x, y-1, z$, and $\mathrm{O}\left(\mathrm{H}_{2} \mathrm{O}\right)-\mathrm{H} \cdots \mathrm{O}(2), \frac{1}{2}+x, \frac{1}{2}-y, \frac{3}{4}-z$ The lengths of these bonds, 2.83 and $2.75 \AA$, indicate that they are fairly weak interactions, and this is confirmed by the significantly higher value of $B$ found for the water oxygen atoms than for the other light atoms.

The remaining hydrogen atom available for bond formation, that attached to $\mathrm{N}(3)$ bonds to $\mathrm{O}(\mathrm{l})$ of a neighbouring molecule, as shown. In turn, the atom $N(3)$ of that molecule is involved in an analogous bond to the $\mathrm{O}(1)$ of the molecule one translation along $a$ from the first, and so on, forming a spiralling system of interactions.

The sole van der Waals contacts between molecules occur at $C(6)$ and $N(3)$ where molecules approach to within $3.76 \AA$ of each other.

We are indebted to the Computational Laboratories of the Massachusetts Institute of Technology and the Applied Physics Laboratory of Johns Hopkins University for the use of their facilities.

We gratefully acknowledge the financial assistance of the National Institutes of Health and the National Science Foundation of the U.S.A.

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# The Crystal Structure of Acrylic Acid 

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(Received 8 June 1962)
The structure of crystalline acrylic acid $\left(\mathrm{C}_{3} \mathrm{H}_{4} \mathrm{O}_{2}\right)$ has been determined at $-135{ }^{\circ} \mathrm{C}$ by X-ray diffraction methods. The crystals were found to be orthorhombic, space group $I b a m-D_{2 h}^{26}$ with 8 molecules per unit cell. The cell constants are

$$
a=9.966 \pm 0.007, b=11.744 \pm 0.013, c=6.306 \pm 0.016 \AA
$$

The structure was refined by least-squares techniques. The molecules are planar hydrogen bonded dimers lying on crystallographic mirror planes.

## Introduction

The structure of acrylic acid, $\left(\mathrm{CH}_{2}=\mathrm{CH}-\mathrm{COOH}\right)$, has been examined as part of a series of investigations of the structures of low molecular weight carboxylic acids. Acrylic acid is the simplest unsaturated aliphatic acid. Its melting point $\left(12{ }^{\circ} \mathrm{C}\right)$ is some $30^{\circ}$ higher than that of propionic acid $\left(-22^{\circ} \mathrm{C}\right)$ and very close to that of acetic acid $\left(16.5{ }^{\circ} \mathrm{C}\right)$. This fact suggested
that the hydrogen bonding system is of the infinite chain type as in acetic acid (Jones \& Templeton, 1958) and not dimeric as in propionic acid (Strieter, Templeton, Scheerman \& Sass, 1962). This investigation shows, however, that acrylic acid crystallizes in hydrogen bonded dimers across a symmetry center, the observed high melting point not being a consequence of the hydrogen bonding.

## Experimental

Crystals of acrylic acid were obtained by partially freezing the pure 'Eastman grade' liquid in a tear bottle, pouring off the liquid phase and then cooling to $-10^{\circ} \mathrm{C}$. Generally, several suitable crystals would be found projecting into the void left by the liquid. These were broken off and shaken onto an aluminum plate cooled by a block of dry ice. The crystals were then mounted in glass capillaries, carried to the X-ray machine on a piece of dry ice and placed in a goniometer head bathed in a stream of cold nitrogen gas. The temperature of the crystals was maintained at $-35{ }^{\circ} \mathrm{C}$ during alignment and was then lowered to $-115 \pm 3{ }^{\circ} \mathrm{C}$ to record the diffraction data. No phase change was observed between the melting point and $-115{ }^{\circ} \mathrm{C}$. All diffraction patterns were obtained with $\mathrm{Cu} K \alpha$ radiation on Kodak no-screen medical X-ray film. Oscillation and Weissenberg photographs show the crystals to be body-centered orthorhombic. The space group is either $\operatorname{Iba} 2\left(C_{2 v}^{21}\right)$ or $\operatorname{Ibam}\left(D_{2 h}^{26}\right)$ derived from the following systematic absences:

$$
\begin{aligned}
& h k l \text { absent if } h+k+l=2 n+1 \\
& 0 k l \text { absent if } k=2 n+1 \\
& h 0 l \text { absent if } h=2 n+1 \\
& (h k 0 \text { absent if } h+k=2 n+1) .
\end{aligned}
$$

The cell constants $a$ and $b$ were determined from a zero layer Weissenberg photograph which was calibrated with a superimposed $\mathrm{NaCl}\left(a_{0}=5 \cdot 637 \AA\right)$ powder pattern. The $c$ axis was determined from a rotation photograph calibrated by a superimposed NaCl rotation diffraction pattern. The following values were obtained:

$$
\begin{gathered}
a=9.966 \pm 0.007, b=11.744 \pm 0.013 \\
c=6.306 \pm 0.016 \AA
\end{gathered}
$$

These constants, assuming eight molecules/unit cell, result in a calculated density for solid acrylic acid of $1.296 \mathrm{~g} . \mathrm{cm}^{-3}$. There has been no experimental density determination made on solid acrylic acid, but density of the liquid at $12{ }^{\circ} \mathrm{C}$ has been reported to be 1.0600 g.cm ${ }^{-3}$ (Moureau \& Boutaric, 1920).

## Treatment of data

Weissenberg photographs of the $l=0,1,2$ and 3 layers were taken with sets of four films and the intensities were estimated visually, a standard intensity strip being employed. The various layers were correlated by comparison with a photograph which contained 15 -minute exposures of a 25 degree portion of each layer. The intensities were corrected with the proper Lorentz, polarization and Tunnell factors. No correction for absorption was deemed necessary since $\mu r$ was calculated to be 0.148 .

If the correct space group is the centric Ibam, the molecules must be in the special positions

$$
8(j):\left(000, \frac{1}{2} \frac{1}{2} \frac{1}{2}\right)+x y 0, \bar{x} \bar{y} 0, \bar{x} y \frac{1}{2}, x \bar{y} \frac{1}{2} .
$$

Under these restrictions, the calculated unitary structure factors of constant $h$ and $k$ values should be equal for all values of $l$ of the same parity. For particular values of $h$ and $k$ then, $F_{h k 0} / \hat{f}_{n k 0}=F_{h k 2} / \hat{f}_{h k i}$ where $\hat{f}$ is the unitary form factor. Taking $\hat{f}=$ $\frac{1}{2}\left(\frac{1}{8} f_{o}+\frac{1}{6} f_{c}\right)$, a plot of $F_{h k 0} / \hat{f}_{h k 0}$ vs. $F_{h k 2} / \hat{f}_{h k 2}$ for 25 random pairs of reflections is shown in Fig. 1. The


Fig. 1. A plot of $F_{h k_{0} \mid}\left(\hat{f_{h k 0}}\right.$ versus $F_{h k_{2}} / \hat{f_{h k_{2}}}$ for 25 pairs
data for this plot are on an arbitrary scale and no correction was made for the effects of thermal vibration. The fact that the plotted points fall close to a straight line influenced us to assume that the centric space group is the correct one.

An attempt was made to determine a correct trial structure from the Patterson projection $P(u v)$. The resulting map was poorly resolved and several attempts failed to lead to a solution. Because the atoms are in special positions, it is possible effectively to subtract out the body centered molecules and reduce the problem to that of a primitive cell. In the space group Ibam with molecules in position $8(j)$, the structure factors of the $l=0$ and $l=1$ layer are respectively
$F_{h k 0}=8 \sum_{j} f_{j}\left(\cos 2 \pi h x_{j}\right)\left(\cos 2 \pi k y_{j}\right) ;(h+k=2 n)$
$F_{h k 1}=-8 \Sigma_{j} f_{j}\left(\sin 2 \pi h x_{j}\right)\left(\sin 2 \pi k y_{j}\right) ;(h+k=2 n+1)$.
In the primitive space group $P b a m$, the special position $4 g$ are related to positions $8(j)$ of Ibam by removing the body centering condition. The corresponding expression for the $l=0$ structure factors are
$F_{h k 0}=4 \sum_{j} f_{j}\left(\cos 2 \pi h x_{j}\right)\left(\cos 2 \pi h y_{j}\right) ;(h+k=2 n)$
$F_{h k 0}=-4 \Sigma_{j} f_{j}\left(\sin 2 \pi h x_{j}\right)\left(\sin 2 \pi k y_{j}\right) ;(h+k=2 n+1)$.
Thus by multiplying the observed values of $F_{h k 1}$ by the ratio of unitary form factors $\hat{f}_{h k 0} / \hat{f}_{h k 1}$, the combined zero and first layer data can be used to calculate the Patterson projection for the primitive cell. It is
obvious that this process is quite closely related to calculating the $w=0$ Patterson section of the bodycentered cell but requires the use of much fewer data. This synthesis was well resolved and led immediately to a trial structure consisting of hydrogen-bonded dimers across a center of symmetry. Only the five 'heavy' atoms were located. The trial structure was refined by a least-squares treatment of the $l=0,1$ and 2 layer data. A total of 119 observed and 51 unobserved (too weak) reflections were used in the calculation. The unobserved reflections were given a value of one-half the minimum observable value. These calculations were carried out on the Rice Computer (Rice University, 1962) with a program written for that machine. The diagonal approximation of the normal equation matrix was used. All nondiagonal elements were set to zero. Several iterations


Fig. 2. $\varrho(x y)_{o}-\varrho(x y)_{c}$ :
(a) centered cell; (b) primitive cell.
were first carried out on only the positional parameters. The thermal vibration parameters were allowed to vary during the final cycles of refinement. After

Table 1. Acrylic acid: Final values of $F_{c}$ and $F_{o}$ including hydrogen atom contributions

convergence had been obtained, several additional cycles were run during which the positional parameters oscillated around the reported values by no more than $\pm 0.0001$ (fractional part of cell dimension) and the thermal vibration parameters oscillated by less than $\pm 0.01 \AA^{2}$. The weighting scheme used was

$$
\begin{aligned}
& w=1 /\left|F_{o}\right| \quad \text { if } \quad\left|F_{o}\right|>4\left|F_{\min }\right| \\
& w=1 / 4\left|F_{\min }\right| \quad \text { if } \quad\left|F_{o}\right|<\left|F_{\min }\right|
\end{aligned}
$$

Individual isotropic temperature factors were refined for all atoms. The carbon and oxygen form factors were those of Berghuis, Haanappel, Potters, Loopstra, MacGillavry \& Veenendaal (1955). The form factors were stored in a table of 64 entries for equally spaced values of $\sin \theta / \lambda$ from 0 to 1 ; linear interpolation was used. During refinement the value of $R$ (observed reflections only) was reduced to $13 \cdot 0$. At this point electron density difference syntheses were calculated to locate the hydrogen atoms. The calculated structure factors used were those of the final cycle of least squares including only the heavy atoms. The projections $\left(\varrho_{o}-\varrho_{c}\right.$ ) down the $c$ axis are shown in Fig. 2 for the body centered and primitive structure. The contours start at 0.2 electrons/ $\AA^{2}$ and are at intervals of 0.2 electrons in the body centered structure and at intervals of $0 \cdot 1$ electrons in the primitive structure. In the body centered cell three of the four hydrogen atoms overlap with equivalent hydrogen atoms from adjacent molecules. In the primitive structure all hydrogen atoms are resolved. The coordinates of the various hydrogen atoms were measured, included in the calculated structure factors and the heavy atom positions were further refined. The final value of $R$ including hydrogen atoms is $10 \cdot 2$. The form factor used by hydrogen was that of McWeeny (1951). The final calculated and observed structure factors are listed in Table 1. (These values are reduced by a factor of 8.) The final values of the atomic parameters are listed in Table 2. The estimated standard deviations reported in this table were obtained from the least-squares treatment.

Table 2. Acrylic acid:

## Final values of atomic parameters

|  | $x$ | $y$ | $z$ | $B$ | $\sigma(x)$ | $\sigma(y)$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{O}(1)$ | 0.1655 | 0.0635 | 0 | $4 \cdot 33$ | 0.0007 | 0.0006 |
| $\mathrm{O}(2)$ | -0.0402 | 0.1362 | 0 | $4 \cdot 42$ | 0.0007 | 0.0006 |
| $\mathrm{C}(1)$ | 0.0854 | 0.1487 | 0 | $3 \cdot 49$ | 0.0009 | 0.0008 |
| $\mathrm{C}(2)$ | 0.1492 | 0.2615 | 0 | $4 \cdot 33$ | 0.0010 | 0.0009 |
| $\mathrm{C}(3)$ | 0.0765 | 0.3538 | 0 | $4 \cdot 63$ | 0.0010 | 0.0009 |
| $\mathrm{H}(1)$ | 0.113 | -0.004 | 0 | $4.5^{*}$ |  |  |
| $\mathrm{H}(2)$ | 0.256 | 0.258 | 0 | $4 \cdot 5^{*}$ |  |  |
| $\mathrm{H}(3)$ | 0.130 | 0.442 | 0 | $4 \cdot 5^{*}$ |  |  |
| $\mathrm{H}(4)$ | -0.014 | 0.347 | 0 | $4 \cdot 5^{*}$ |  |  |
| * Assumed values. |  |  |  |  |  |  |

## Discussion of the structure

Acrylic acid crystallizes as planar, hydrogen bonded dimers. The arrangement of a single layer of molecules


Fig. 3. Molecular configuration and dimensions in the $z=0$ plane.
in the $z=0$ plane is shown in Fig. 3. Bond distances and angles are presented in Table 3. Within experimental error, the $\mathrm{C}-\mathrm{C}$ double bond is of expected length and the $\mathrm{C}-\mathrm{C}$ single-bond length is that of a normal $s p^{2}-s p^{2}$ bond. The $\mathrm{C}=\mathrm{O}$ bond length is somewhat longer and the $\mathrm{C}-\mathrm{O}$ bond length is shorter

Table 3. Acrylic acid: Bond lengths and angles

| Bond | Length $\pm$ e.s.d. | Angle |  |
| :--- | :--- | :--- | :--- |
|  |  |  |  |
| $\mathrm{C}(1)-\mathrm{C}(2)$ | $1 \cdot 47 \pm 0 \cdot 015 \AA$ | $\mathrm{C}(2)-\mathrm{C}(1)-\mathrm{O}(2)$ | $122 \pm 1^{\circ}$ |
| $\mathrm{C}(2)-\mathrm{C}(3)$ | $1 \cdot 30 \pm 0 \cdot 015$ | $\mathrm{C}(2)-\mathrm{C}(1)-\mathrm{O}(1)$ | $116 \pm 1$ |
| $\mathrm{C}(1)-\mathrm{O}(1)$ | $1 \cdot 28 \pm 0 \cdot 014$ | $\mathrm{O}(1)-\mathrm{C}(1)-\mathrm{O}(2)$ | $122 \pm 1$ |
| $\mathrm{C}(1)-\mathrm{O}(2)$ | $1 \cdot 26 \pm 0 \cdot 014$ | $\mathrm{C}(1)-\mathrm{O}(1)-\mathrm{H}(1)$ | 111 |
| $\mathrm{O}(1)-\mathrm{H}(1)$ | $0 \cdot 95 \pm 0 \cdot 1$ | $\mathrm{C}(1)-\mathrm{O}(1) \cdots \mathrm{O}(2)$ | $114 \pm 1$ |
| $\mathrm{C}(2)-\mathrm{H}(2)$ | $1 \cdot 06 \pm 0 \cdot 1$ | $\mathrm{C}(1)-\mathrm{O}(2) \cdots \mathrm{O}\left(11^{\prime}\right) 125 \pm 1$ |  |
| $\mathrm{C}(3)-\mathrm{H}(3)$ | $1 \cdot 16 \pm 0 \cdot 1$ | $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{H}(2)$ | 114 |
| $\mathrm{C}(3)-\mathrm{H}(4)$ | $0 \cdot 91 \pm 0 \cdot 1$ | $\mathrm{C}(3)-\mathrm{C}(2)-\mathrm{H}(2)$ | 126 |
| $\mathrm{O}(1) \cdots \mathrm{O}\left(2^{\prime}\right)$ | $2 \cdot 66 \pm 0 \cdot 014$ | $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{C}(3)$ | $120 \pm 1$ |
|  |  | $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{H}(4)$ | 129 |
|  |  | $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{H}(3)$ | 119 |
|  |  | $\mathrm{H}(3)-\mathrm{C}(3)-\mathrm{H}(4)$ | 112 |

than the corresponding values observed in normal saturated carboxylic acids, but they agree well with observed distances in unsaturated and aromatic acids. All bond angles are normal. A representative listing of carboxylic acids is presented in Table 4. With the exception of 2 -naphthoic acid, the carboxyl group dimensions fall into two classes. In the saturated dimeric acids the two $\mathrm{C}-\mathrm{O}$ bonds lengths differ by approximately $0 \cdot 11 \AA$. These two distances in the unsaturated acids are, however, very nearly the same, differing by only $0.03 \AA$ on the average. It is of interest that although the dimeric gas phase formic and acetic acid dimensions are consistent with crystalline dimeric propionic, butyric and valeric acids, the crystalline forms of formic and acetic acids, which consist of infinite hydrogen bonded chains,

Table 4. Bond distances and angles in various carboxylic acids

| Compound | $\mathrm{C}-\mathrm{OH}$ | $\mathrm{C}=\mathrm{O}$ | $\mathrm{R}-\mathrm{C}-\mathrm{OH}$ | $\mathrm{R}-\mathrm{C}-\mathrm{O}$ | $\mathrm{O}=\mathrm{C}-\mathrm{OH}$ | $\mathrm{O} \cdots \mathrm{O}$ | Compound type | Reference |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Formic acid | $1.36 \AA$ | $1 \cdot 25 \AA$ |  |  | $121^{\circ}$ | $2.73 \AA$ | Gas, dimer | Karle \& Brockway <br> (1944) |
| Acetic acid | $1 \cdot 36$ | 1.25 | $110^{\circ}$ | $120^{\circ}$ | 130 | $2 \cdot 76$ | Gas, dimer | Karle \& Brockway (1944) |
| Propionic acid | $1 \cdot 33$ | $1 \cdot 22$ | 113 | 125 | 122 | $2 \cdot 64$ | Crystal, dimer | Strieter, Templeton, Scheuerman \& Sass (1962) |
| Butyric acid | $1 \cdot 34$ | 1.22 | 1145 | 126 | 1195 | $2 \cdot 66$ | Crystal, dimer | Strieter \& Templeton (1962) |
| Valeric acid | $1 \cdot 35$ | $1 \cdot 26$ | 1165 | 125 | 1185 | $2 \cdot 63$ | Crystal, dimer | Scheuerman \& Sass (1962) |
| Average | $1.35 \AA$ | $1 \cdot 24 \AA$ | $114^{\circ}$ | $124^{\circ}$ | $122^{\circ}$ |  |  |  |
| Acetylene dicarboxylic acid | $1.27 \AA$ | $1.26 \AA$ | $117^{\circ}$ | $123^{\circ}$ | $120^{\circ}$ |  | Crystal, dimer | Dumitz \& Robertson (1947) |
| Acrylic acid | $1 \cdot 28$ | 1.26 | 116 | 122 | 122 | $2 \cdot 66$ | Crystal, dimer | This paper |
| Benzoic acid | $1 \cdot 29$ | $1 \cdot 24$ | 118 | 122 | 122 | $2 \cdot 64$ | Crystal, dimer | Sim, Robertson \& Goodwin (1955) |
| Formic acid | $1 \cdot 26$ | $1 \cdot 23$ |  | 123 |  | $2 \cdot 58$ | Crystal, infinite chain | Holtzberg, Post \& Fankucken (1953) |
| Acetic acid | $1 \cdot 29$ | $1 \cdot 24$ | 116 | 122 | 122 | $2 \cdot 61$ | Crystal, infinite chain | Jones \& Templeton (1958) |
| 1-Naphthoic acid | $1 \cdot 28$ | 1.25 | 122 | 127 | 110 | 2.58 | Crystal, dimer | Trotter (1960) |
| 2-Naphthoic acid | $1 \cdot 37$ | $1 \cdot 33$ | 127 | 122 | 112 | 2.54 | Crystal, dimer | Trotter (1961) |

contain $\mathrm{C}-\mathrm{O}$ distances very similar to those of the unsaturated acids.

The various bonded hydrogen distances are probably accurate to only about $\pm 0 \cdot 1 \AA$. The average $\mathrm{C}-\mathrm{H}$ distance of $1.04 \AA$ is fortuitously close to the accepted value of $1.07 \AA$. The measured value of $0.95 \AA$ for the $\mathrm{O}-\mathrm{H}$ distance leaves no doubt that the hydrogen atom is singly bonded and not symmetrical in the hydrogen bond as might be implied from the similarity of the two C-O distances.

All heavy atom intramolecular contact distances less than or equal to $3 \cdot 50 \AA$ are, in the same plane

$$
\begin{array}{ll}
\mathrm{O}(1)-\mathrm{O}(2) & 2 \cdot 66 \AA \text { (hydrogen bond) } \\
\mathrm{C}(2)-\mathrm{O}(2) & 3 \cdot 32 \AA
\end{array}
$$

between planes

$$
\begin{array}{cl}
\mathrm{C}(3)-\mathrm{C}\left(3^{\prime \prime}\right) & 3 \cdot 50 \AA \\
\mathrm{C}(1)-\mathrm{O}\left(2^{\prime \prime}\right) & 3 \cdot 19 \AA \\
\mathrm{O}(1)-\mathrm{O}\left(2^{\prime \prime}\right) & 3 \cdot 49 \AA
\end{array}
$$

Note added October 23, 1962: It has recently come to the attention of the authors that a paper discussing the structure of acrylic acid was presented by Isamu Nitta at the Cambridge meeting of the International Union of Crystallography in 1960. The paper is abstracted in Acta Cryst. (1960), 13, 1035. The gross structure there presented is the same as found in the present paper, but the bond distances reported for the carboxyl group are significantly different. So far as the authors know, no further account of this work has been published.

These studies were supported by a grant from the Robert A. Welch Foundation of Texas. The use of the Rice Computer was made possible by the Atomic Energy Commission, contract number AT-(40-1)-1825.

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